

## LECTURE 11:

# BIFURCATIONS IN 2-D MAPS

We prelude our discussion of 2-D nonlinear maps with a discussion of linear 2-D maps.

### 3.1 Linear 2-D Maps

Linear maps on  $\mathbb{R}^2$  are linear transformations defined by  $2 \times 2$  matrices ie.  $\underline{v} \mapsto A\underline{v}$ , where  $A$  is a  $2 \times 2$  matrix with

$$A(a\underline{v} + b\underline{w}) = aA\underline{v} + bA\underline{w},$$

for  $a, b \in \mathbb{R}$ ,  $\underline{v}, \underline{w} \in \mathbb{R}^2$  and

$$A\underline{v} = A \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}.$$

Every linear map has a fixed point at the origin.

*Proof.* Exercise. □

Recall:  $\lambda$  is an eigenvalue of  $A$  if  $\exists \underline{v} \neq 0$  such that  $A\underline{v} = \lambda\underline{v}$ . Suppose  $\underline{v}_0$  is an eigenvector with eigenvalue  $\lambda$ . Then  $A\underline{v}_0 = \lambda\underline{v}_0$ .

Applying  $A$  times gives

$$\underline{v}_1 = A\underline{v}_0 = \lambda\underline{v}_0$$

$$\underline{v}_2 = A\underline{v}_1 = A\lambda\underline{v}_0 = \lambda A\underline{v}_0 = \lambda^2\underline{v}_0.$$

$\therefore$

$$\underline{v}_n = \lambda^n \underline{v}_0 \quad \text{and} \quad \underline{v}_{n+1} = A\underline{v}_n,$$

so that it behaves like a 1-D map under iteration.

Example Suppose  $A \in \mathbb{R}^2$  is defined by

$$A \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & 0 \\ 0 & b \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} ax \\ by \end{bmatrix}.$$

Then

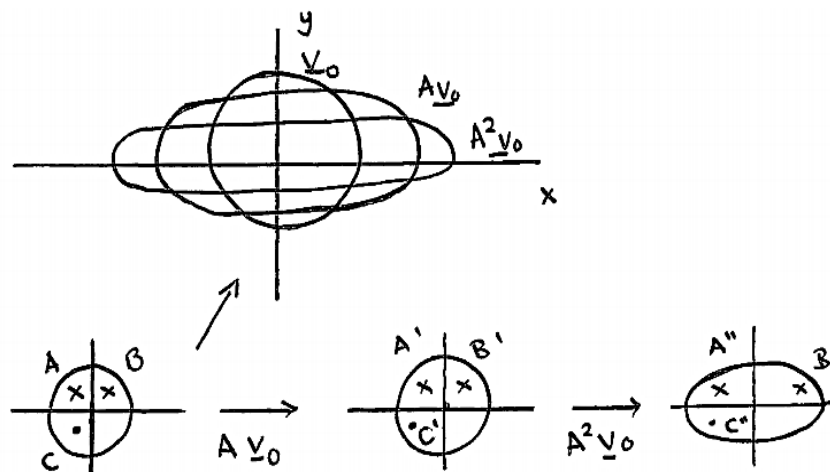
$$A^n = \begin{bmatrix} a^n & 0 \\ 0 & b^n \end{bmatrix}$$

so that

$$A^n \begin{bmatrix} x_0 \\ y_0 \end{bmatrix} = \begin{bmatrix} a^n x_0 \\ b^n y_0 \end{bmatrix} = \begin{bmatrix} x \\ y \end{bmatrix}.$$

It follows that the unit disc is mapped into an ellipse with axes  $|a|^n$  and  $|b|^n$  along  $x$  and  $y$  respectively.

- If both  $|a| < 1$  and  $|b| < 1$ , the axes shrink as  $n \rightarrow \infty$ , and we have a sink.
- If  $|a| > 1 > |b|$ , the  $x$ -axis ( $|a|^n$ ) grows, while the  $y$ -axis ( $|b|^n$ ) shrinks as  $n \rightarrow \infty$ , and we have a saddle.
- If  $|a| > 1$  and  $|b| > 1$ , both axes expand as  $n \rightarrow \infty$ , and we have a source.



The determinant  $\det A$  measures the area transformed by  $A$ .

**Definition 6.** If  $\det A = \pm 1$ , we have an area-preserving map.

**Definition 7.**  $A$  is a hyperbolic map if no eigenvalues have magnitude one.

**Lemma 2.** If  $A(\underline{y})$  is a linear map on  $\mathbb{R}^2$ , then

1. The origin is a sink if all eigenvalues of  $A$  have magnitude  $< 1$ .
2. The origin is a source if all eigenvalues of  $A$  have magnitude  $> 1$ .
3. If a hyperbolic map has eigenvalues  $> 1$  and  $< 1$ , the origin is a saddle.

### 3.2 Nonlinear Maps and Jacobians

Consider  $\underline{x}_{n+1} = F(\underline{x}_n)$  with fixed pt  $\underline{X} = F(\underline{X})$ .

Perturb away from  $\underline{X}$ :  $\underline{\xi}_n = \underline{x}_n - \underline{X}$ .

Then

$$\underline{x}_{n+1} = \underline{\xi}_{n+1} + \underline{X} = F(\underline{\xi}_n + \underline{X}) = F(\underline{X}) + DF \cdot \underline{\xi}_n + o(|\underline{\xi}_n|),$$

as  $\underline{\xi}_n \rightarrow 0$ , where

$$DF(\underline{X}) = [\nabla F]_{\underline{X}} = \left[ \frac{\partial f_i}{\partial x_j} \right]_{\underline{X}} = J$$

is the Jacobian matrix of  $F$ , evaluated at  $\underline{X}$ .

Thus

$$\underline{\xi}_{n+1} = J \underline{\xi}_n$$

and so stability reduces to determining the nature of fixed points of a linearised 2-D map,  $J$ . This result readily extends to  $n$ -D.

### 3.3 Periodic Orbits

Suppose  $A$  has a periodic orbit  $\{\underline{X}_1, \dots, \underline{X}_p\} = \underline{S}$  of period  $p$ . Then  $\underline{X}_{r+1} = \underline{F}(\underline{X}_r)$  for  $r = 1, \dots, p-1, p$  with  $\underline{X}_1 = \underline{F}(\underline{X}_p)$  where

$$\underline{F}(\underline{S}) = \{\underline{F}(\underline{X}_1), \dots, \underline{F}(\underline{X}_p)\} = \{\underline{X}_2, \dots, \underline{X}_1\}.$$

Thus  $\underline{S}$  is an invariant set under  $\underline{F}$ .

To consider the stability of a  $p$ -cycle, we consider the  $p^{\text{th}}$  iteration of the map.

### 3.4 2-Cycles

Consider the 2-cycle  $\{\underline{X}, \underline{Y}\}$  such that  $\underline{Y} = \underline{F}(\underline{X})$ ,  $\underline{X} = \underline{F}(\underline{Y})$ . Then

$$\underline{G}(\underline{X}) = \underline{F}^2(\underline{X}) = \underline{F}(\underline{F}(\underline{X})).$$

A fixed point of  $\underline{G}(\underline{X}) = \underline{X} = \underline{F}(\underline{F}(\underline{X}))$  is either a fixed point of  $\underline{F}$  or a period-2 cycle.

Proof: Exercise

Suppose  $\underline{F}$  is twice continuously differentiable. Then, linearising  $\underline{G}$  about  $\underline{X}$ :

$$\underline{\xi}_{n+1} = \left[ \frac{\partial G_i}{\partial x_j} \right]_{\underline{X}} \underline{\xi}_n = \underline{K}(\underline{X}) \underline{\xi}_n.$$

However stability is given by the chain rule:

$$\underline{K}(\underline{X}) = \underline{D}\underline{G} = \underline{D}(\underline{F}(\underline{F}(\underline{X}))) = \underline{D}\underline{F}(\underline{F}(\underline{X}))\underline{D}\underline{F}(\underline{X}) = \underline{D}\underline{F}(\underline{Y})\underline{D}\underline{F}(\underline{X}).$$

Therefore

$$\underline{K}(\underline{X}) = \underline{J}(\underline{F}(\underline{X}))\underline{J}(\underline{X}) = \underline{J}(\underline{Y})\underline{J}(\underline{X}).$$

Thus for 2-D maps, the stability of  $\underline{X}$  and  $\underline{Y}$  is given by the eigenvalues of the product of two Jacobians,  $\underline{D}\underline{F}(\underline{X})$  and  $\underline{D}\underline{F}(\underline{Y})$ , evaluated at  $\underline{X}$  and  $\underline{Y}$ .

If  $f$  is a 1-1 map, its inverse  $f^{-1}$  is a function and is defined.

### 3.5 Lyapunov Exponents and Numbers

Above we saw that perturbations  $\xi_n = x_n - x_*$  from the fixed point  $x_*$  of a 1-D map satisfy, under iteration on  $n$

$$\xi_n = |f'(x_*)|^n \xi_0.$$

Thus, when  $f'(x_*) > 1$ , the orbit of each point close to  $x_*$  diverges at a multiplicative rate of approximately  $f'(x_*)$ .

For a periodic-cycle of period  $K$ , stability is given by  $f'(x_1)f'(x_2)\dots f'(x_K) = A$ . If  $A > 1$ , the orbit or each neighbour  $x$  of the period  $K$ -cycle diverges at a rate of  $\approx A$  after  $K$  iterations. The average multiplicative rate of separation is  $A^{1/k}$  per iterate.

**Definition 8.** The Lyapunov number (L.N.) quantifies this average rate of separation of its very close to  $x_*$ .

**Definition 9.** The Lyapunov exponent (L.E.) is the natural logarithm of the Lyapunov number.

Examples

1. A L.N. of 2 (ie. L.E  $\ln 2$ ) for the orbit of  $x_*$  means the distance between the orbit of  $x_*$  means the distance between the orbit of  $x_*$  and of  $x_* + \xi$  for  $|\xi| \ll 1$ , will double after each iteration.
2. A L.N. of 1/2, means the instance halves at each iteration and the orbits of  $x_*$  and  $x_* + \xi$  move rapidly together.
3. A L.N. of 1 means no change: its LE= 0.

L.N.s and L.E.s are important when applied to chaotic (non periodic) orbits, where there is sensitive dependence on i.c.

**Definition 10.** Let  $f$  be a smooth map  $\mathbb{R} \rightarrow \mathbb{R}$ . The L.N.  $L(x_1)$  of the orbit  $\{x_1, x_2, \dots\}$  is defined as

$$LN(x_1) = \lim_{n \rightarrow \infty} (|f'(x_1)||f'(x_2)| \dots |f'(x_n)|)^{1/n},$$

if this limit exists and

$$L.E.(x_1) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{m=1}^n \ln |f'(x_m)|.$$

Since  $\ln L = L.E.$ , the L.E. exists if LN exists and is non-zero.

N.B.

1. The L.N. of a fixed pt  $x_*$  for a 1-D map is  $|f'(x_*)|$  and its L.E. is  $\ln f'(x_1)$ .
2. The L.E. of a period K-cycle is

$$L.E.(x_1) = \frac{1}{K} \sum_{m=1}^K \ln |f'(x_m)|,$$

and the L.N.,  $L(x_1) = e^{L.E.(x_1)}$  describes the average local stretching per iterate.

3. A chaotic orbit has  $L.E. > 0$ .

Eg.: For the Bernoulli shift map  $f(y) = 2y \bmod 1$ , ( $y_{n+1} = 2y_n \bmod 1$ ), the  $L.E.$  is

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{m=1}^n \ln |f'(y_m)| = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{m=1}^n \ln 2 = \ln 2.$$