

# FOURIER SERIES AND PDEs

## LECTURE 6

### INITIAL AND BOUNDARY VALUE PROBLEM (IBVP)

We now consider finding the solution of the heat equation

$$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.31)$$

subject to the initial condition

$$T(x, 0) = f(x), \quad 0 \leq x \leq L, \quad (3.32)$$

and the boundary conditions

$$T(0, t) = 0 \text{ and } T(L, t) = 0 \text{ for } t > 0. \quad (3.33)$$

In view of our preceding discussion we look for a solution as an infinite sum

$$T(x, t) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right) e^{-n^2\pi^2\kappa t/L^2}. \quad (3.34)$$

**Example 3.1** Solve the IBVP for the case

$$T(x, 0) = \sin\left(\frac{\pi x}{L}\right) + \frac{1}{2} \sin\left(\frac{2\pi x}{L}\right) = f(x), \quad 0 \leq x \leq L. \quad (3.35)$$

From equation (3.34) we see that

$$T(x, 0) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right). \quad (3.36)$$

Comparing terms we see that  $a_1 = 1$ ,  $a_2 = 1/2$  and  $a_n = 0$  ( $n \geq 3$ ) so that the solution is

$$T(x, t) = \sin\left(\frac{\pi x}{L}\right) e^{-\pi^2\kappa t/L^2} + \frac{1}{2} \sin\left(\frac{2\pi x}{L}\right) e^{-4\pi^2\kappa t/L^2}. \quad (3.37)$$

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### Application of Fourier series

To solve for more general initial conditions, we can use Fourier series to determine the constants  $a_n$ :

$$T(x, 0) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right), \quad 0 \leq x \leq L. \quad (3.38)$$

The question is now, given  $f(x)$ , can it be expanded as a Fourier sine series

$$f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right), \quad 0 \leq x \leq L? \quad (3.39)$$

From the lectures on Fourier series, we know that such an expansion exists if *e.g.*  $f$  is piecewise continuously differentiable on  $[0, L]$ . The coefficients  $a_n$  are determined by the orthogonality relations:

$$\int_0^L \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx = \begin{cases} 0, & m \neq n, \\ \frac{1}{2}L, & m = n. \end{cases} \quad (3.40)$$

Thus

$$a_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx. \quad (3.41)$$

**Example 3.2** Find the solution of the IBVP when

$$f(x) = \begin{cases} 0 & \text{for } 0 \leq x \leq L_1 \text{ and } L_2 \leq x \leq L, \\ 1 & \text{for } L_1 < x < L_2. \end{cases} \quad (3.42)$$

Here  $f(x)$  has the Fourier sine expansion

$$\frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} \left[ \cos\left(\frac{n\pi L_1}{L}\right) - \cos\left(\frac{n\pi L_2}{L}\right) \right] \sin\left(\frac{n\pi x}{L}\right), \quad (3.43)$$

and the solution of IBVP is

$$T(x, t) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} \left[ \cos\left(\frac{n\pi L_1}{L}\right) - \cos\left(\frac{n\pi L_2}{L}\right) \right] \sin\left(\frac{n\pi x}{L}\right) e^{-n^2 \pi^2 \kappa t / L^2}. \quad (3.44)$$

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### Uniqueness

We have constructed a solution of our IBVP, and found a formula for it as the sum of an infinite series, but is it the only solution?

**Theorem 3.1** The IBVP has only one solution.

*Proof.* Let  $U$  be a solution of the same IBVP, *i.e.*

$$\frac{\partial U}{\partial t} = \kappa \frac{\partial^2 U}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.45)$$

subject to the initial condition

$$U(x, 0) = f(x), \quad 0 \leq x \leq L, \quad (3.46)$$

and the boundary conditions

$$U(0, t) = 0 \text{ and } U(L, t) = 0 \text{ for } t > 0. \quad (3.47)$$

Now consider the difference  $W := U - T$ . Then  $W$  satisfies the IBVP

$$\frac{\partial W}{\partial t} = \kappa \frac{\partial^2 W}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.48)$$

$$W(x, 0) = 0, \quad 0 \leq x \leq L, \quad (3.49)$$

and the boundary conditions

$$W(0, t) = 0 \text{ and } W(L, t) = 0 \text{ for } t > 0. \quad (3.50)$$

Let

$$I(t) := \frac{1}{2} \int_0^L [W(x, t)]^2 dx. \quad (3.51)$$

Evidently  $I(t) \geq 0$  and  $I(0) = 0$ . By Leibniz's rule,

$$I'(t) = \int_0^L W \frac{\partial W}{\partial t} dx, \quad (3.52)$$

$$= \kappa \int_0^L W \frac{\partial^2 W}{\partial x^2} dx, \quad (3.53)$$

$$= \kappa \int_0^L \left[ \frac{\partial}{\partial x} \left( W \frac{\partial W}{\partial x} \right) - \left( \frac{\partial W}{\partial x} \right)^2 \right] dx. \quad (3.54)$$

On carrying out the integration and using the boundary conditions at  $x = 0$  and  $x = L$  we see that

$$I'(t) = -\kappa \int_0^L \left( \frac{\partial W}{\partial x} \right)^2 dx \leq 0, \quad (3.55)$$

and, therefore,  $I$  cannot increase. Hence

$$0 \leq I(t) \leq I(0) = 0, \quad (3.56)$$

and  $I(t) = 0$  for every  $t \geq 0$ . Thus

$$\int_0^L [W(x, t)]^2 dx = 0, \quad (3.57)$$

for every  $t \geq 0$  and so  $W = 0$  and  $U = T$ , which proves the theorem.  $\square$

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## Non-zero steady state

It may be that the temperatures of the ends  $x = 0$  and  $x = L$  are prescribed and constant but not equal to zero.

**Example 3.3** Solve the IBVP

$$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.58)$$

subject to the initial condition

$$T(x, 0) = 0, \quad 0 \leq x \leq L, \quad (3.59)$$

and the boundary conditions

$$T(0, t) = T_0 \text{ and } T(L, t) = T_1 \text{ for } t > 0. \quad (3.60)$$

We cannot use separation of variables and Fourier series right at the outset. However, we conjecture that, as  $t \rightarrow \infty$ ,  $T(x, t) \rightarrow U(x)$ , where

$$\kappa \frac{d^2 U}{dx^2} = 0, \quad U(0) = T_0 \text{ and } U(L) = T_1, \quad (3.61)$$

*i.e.*

$$U(x) = T_0 \left(1 - \frac{x}{L}\right) + T_1 \left(\frac{x}{L}\right). \quad (3.62)$$

If we now put  $S(x, t) := T(x, t) - U(x)$ , we find that  $S$  is a solution of the IBVP

$$\frac{\partial S}{\partial t} = \kappa \frac{\partial^2 S}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.63)$$

with

$$S(0, t) = 0 \text{ and } S(L, t) = 0 \text{ for } t > 0, \quad (3.64)$$

and

$$S(x, 0) = -T_0 \left(1 - \frac{x}{L}\right) - T_1 \left(\frac{x}{L}\right). \quad (3.65)$$

In view of the form of the boundary conditions, this IBVP can be solved by our previous methods. The solution is

$$S(x, t) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} [-T_0 + (-1)^n T_1] \sin\left(\frac{n\pi x}{L}\right) e^{-n^2 \pi^2 \kappa t / L^2}, \quad (3.66)$$

and so

$$T(x, t) = T_0 \left(1 - \frac{x}{L}\right) + T_1 \left(\frac{x}{L}\right) + \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} [-T_0 + (-1)^n T_1] \sin\left(\frac{n\pi x}{L}\right) e^{-n^2 \pi^2 \kappa t / L^2}. \quad (3.67)$$

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### Other boundary conditions

Other boundary conditions are possible, *e.g.* at an end which is *thermally insulated* the heat flux is zero. Thus  $-kT_x = 0$  there and, therefore,  $T_x = 0$ . If both ends are thermally insulated we look for separable solutions of the heat equation of the form

$$T(x, t) = F(x)G(t), \quad (3.68)$$

where  $F'(0) = F'(L) = 0$ . We find that  $F'' = -\lambda^2 F$ ,  $G' = -\lambda^2 \kappa G$ , and  $F = a \cos(\lambda x)$ , where  $\sin(\lambda L) = 0$  and so  $L$  is one of the numbers  $\{n\pi/L : n = 0, 1, 2, 3, \dots\}$ . The separable solutions in these circumstances are

$$a_0 \text{ and } a_n \cos\left(\frac{n\pi x}{L}\right) e^{-n^2 \pi^2 \kappa t / L^2} \quad (n = 1, 2, 3, \dots). \quad (3.69)$$

Thus if we consider the IBVP

$$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial x^2}, \quad 0 < x < L, t > 0, \quad (3.70)$$

with boundary conditions

$$\frac{\partial T}{\partial x}(0, t) = 0 \text{ and } \frac{\partial T}{\partial x}(L, t) = 0 \text{ for } t > 0, \quad (3.71)$$

and initial condition

$$T(x, 0) = f(x) \text{ for } 0 \leq x \leq L, \quad (3.72)$$

we look for a solution

$$T(x, t) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right) e^{-n^2 \pi^2 \kappa t / L^2}, \quad (3.73)$$

where the prescribed  $f(x)$  has the Fourier cosine expansion

$$f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right), \quad 0 \leq x \leq L. \quad (3.74)$$

The required coefficients are

$$a_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx \quad (n = 0, 1, 2, 3, \dots). \quad (3.75)$$

Note that, as  $t \rightarrow \infty$ ,

$$T(x, t) \rightarrow \frac{1}{2}a_0 = \frac{1}{L} \int_0^L f(s) ds, \quad (3.76)$$

the extreme right-hand side being the mean initial temperature. The uniqueness of  $T(x, t)$ , for a given  $f(x)$ , can be established much as before.